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HIGH RESOLUTION TRANSMISSION MEASUREMENTS OF THE ATMOSPHERE IN --ETC(U)

APR 78 R ZANDER, M V MIGEOTTE, L DELBOUILLE

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INTERIM SCIENTIFIC REPORT

Period : 1 March 1977 - 28 February 1978

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HIGH RESOLUTION TRANSMISSION MEASUREMENTS
OF THE ATMOSPHERE IN THE INFRARED

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Period : 1 March 1977 - 28 February 1978

HIGH RESOLUTION TRANSMISSION MEASUREMENTS
OF THE ATMOSPHERE IN THE INFRARED

prepared by R. ZANDER

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TABLE OF CONTENTS

| | <u>Page</u> |
|--------------------------|-------------|
| I. - INTRODUCTION | 1 |
| II. - SCIENTIFIC RESULTS | 3 |
| III. - ACKNOWLEDGMENTS | 7 |
| IV. - REFERENCES | 8 |
| - APPENDIX A | |
| - APPENDIX B | |

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I. - INTRODUCTION

During the period covered by this report, we have undertaken a certain number of improvements or modifications to the various spectrometric installations, currently operated to obtain very high resolution solar spectra in the infrared. Such spectra have been analyzed in order to, either deduce transmission characteristics of the atmosphere, or to establish the presence and/or the temporal variation of some minor telluric constituents likely to play a role in relation to the problem of the ozone shield depletion by anthropogenic sources.

These experimental efforts concern :

- a. - the stratospheric instrumentation which had suffered important damages on its last 1976 "splash-down" in a lake ; this equipment has been refurbished and optimized for two balloon flights to be carried out at the National Scientific Balloon Facility, Palestine, Texas, in August and September 1978. Initially planned for the Spring of 1978, this campaign had to be delayed, due to late delivery of the optical mirrors which C. rubb-Parsons Ltd., Newcastle upon Tyne, C.B., had to repair (new cerium coatings, polishing and aluminizing). Very high sensitive indium-antimonide (InSb) detectors, operating in the photo-voltaic mode at liquid nitrogen temperature, will be flown and allow observations in the 2 to 5.2 microns range, with a spectral resolution of 0.025 cm^{-1} ;
- b. - the 7.5 meters focal length prism-grating spectrometer of the Jungfrau-Joch which, aside of being used to further record the solar spectrum in the λ 3000 to λ 10,000 Å region, has also been equipped with an InSb detector to record various series of spectra near 2.5 and 3.3 microns, where, respectively, hydrofluoric (HF) and hydrochloric acid (HCl) have their strongest absorption lines. Provisions have also been made (new Bausch and Lomb grating and mercury-cadmium-telluride (Hg-Cd-Te) detectors) in order to be able to operate this instrument, as soon as possible, in the 8 to 13 microns atmospheric window ;

- c. - the one-meter optical path difference Fourier transform interferometer installed at the Jungfraujoch which, till now, has been used to record various spectral intervals between 1.1 and 2.7 microns; it is being modified to increase its scanning speed and to be operational over the whole "InSb" spectral range (1.0 to 5.2 microns). New beam-splitters and Hg-Cd-Te detectors may become available in a near future, in which case this 0.01 cm^{-1} spectral resolution interferometer could also be used out to about 12-13 microns.
-

II. - SCIENTIFIC RESULTS

As will appear throughout this section, a significant effort has been made in order to deduce, as quickly as possible, an ensemble of informations likely to be of interest to that part of the scientific community seeking precise indications concerning the implications of the presence of chloro-fluoro-carbons (CFC's) in the stratosphere. In particular, we have been funded by the Manufacturing Chemists Association (MCA) to establish the feasibility and usefulness of ground monitorings concerning HF, HCl, also methyl-chloride (CH_3Cl), etc....

Other studies based on our infrared solar observations have also been undertaken or pursued.

1. - Balloon observations

- a. - Some earlier balloon observations have been reanalyzed in order to ascertain previous deductions concerning the atmospheric HF and HCl mixing ratio above 30 km altitude; such informations are of interest to the environmentalists, as they provide cumulative indications on the magnitude of the photolysis of the CFC's in the upper stratosphere.

We have deduced the following average mixing ratios above float altitude.

For HF :

| | | | |
|------------|-----------------------|----------------------------------|------|
| April 71 | average above 27.8 km | : $(3.2 \pm .7) \times 10^{-10}$ | ppv. |
| October 74 | " " 27.4 km | : $(4.8 \pm .8)$ | " " |
| April 76 | " " 27.9 km | : $(3.6 \pm .6)$ | " " |

For HCl :

| | | | |
|------------|-----------------------|---------------------------------|------|
| October 74 | average above 30.5 km | : $(1.8 \pm .4) \times 10^{-9}$ | ppv. |
| May 76 | " " 29.7 km | : $(2.4 \pm .5)$ | " " |

These results lead to HF/HCl ratios equal to

$$\begin{aligned}
 & (.29 \pm .11) \quad \text{for October, 1974} \\
 \text{and} \quad & (.16 \pm .06) \quad \text{for April-May, 1976.}
 \end{aligned}$$

HF and HCl will again be reinvestigated on our forthcoming balloon flights.

b. - In our previous report, we had mentioned the significant differences in the concentrations of both CH_4 and N_2O , as deduced from the balloon observations of 1974 and 1976. Using the AFGL-line parameters compilation by McClatchey et al., [1], we have computed synthetic spectra in the $(1260 \pm 5) \text{ cm}^{-1}$ range and compared them to observed ones; a sample comparison is given in Figure 1. The CH_4 and N_2O distributions which best reproduce our observations are given in Figure 2. No special emphasis was given yet to the "reproduction" of the H_2O line at 1260.370 cm^{-1} , as this is mostly produced by moisture inside of the balloon gondola.

2. - Jungfraujoch observations

a. - The double-pass prism-grating spectrometer was run by Professor L. Delbouille and Dr. G. Roland to obtain various sets of spectra near the HF-K(1) and HCl-K(1) lines (other HF and HCl lines were also observed); typical spectra are reproduced on Figures 3 and 4 (see also Ref. [2]).

We have analyzed these sets of observations to deduce average stratospheric HF mixing ratios and column densities of HCl above the Jungfraujoch. The results are summarized in the following tables, which also include the best estimated values, deduced from observations made at other mountain stations.

For HF :

| | | | | |
|--------------|--------------|--|---|-----------------------------|
| St-Michel | (July 66) | average above 20 km : $(1.1 \pm .4) \times 10^{-10} \text{ ppv}$ | | |
| Kitt Peak | (Winter 69) | " | " | $(1.3 \pm .3)$ |
| | (May 77) | " | " | $(2.1 \pm .3)$ |
| Jungfraujoch | (October 76) | " | " | $(2.2 \pm .3)$ |
| | (March 77) | " | " | $(2.5 \pm .3)$ |
| | " | " | " | $(4.6 \pm .5) \text{ Max.}$ |
| | (July 77) | " | " | $(2.8 \pm .3)$ |
| | (October 77) | " | " | $(2.8 \pm .3)$ |
| | (March 78) | " | " | $(3.1 \pm .4)$ |
| | " | " | " | $(4.6 \pm .4) \text{ Max.}$ |

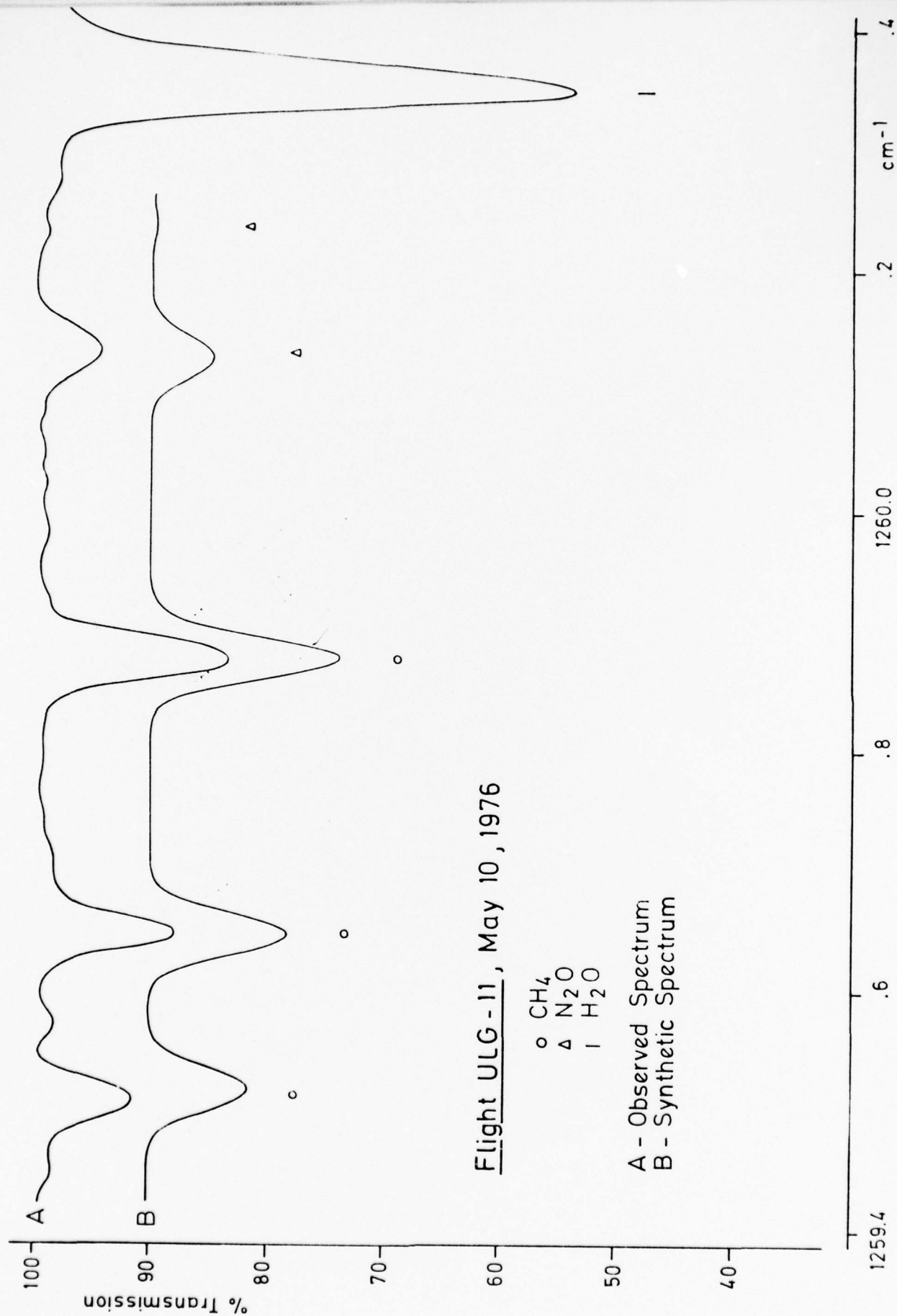


FIGURE 1

MOST REALISTIC PROFILES FITTED TO OBSERVATIONS

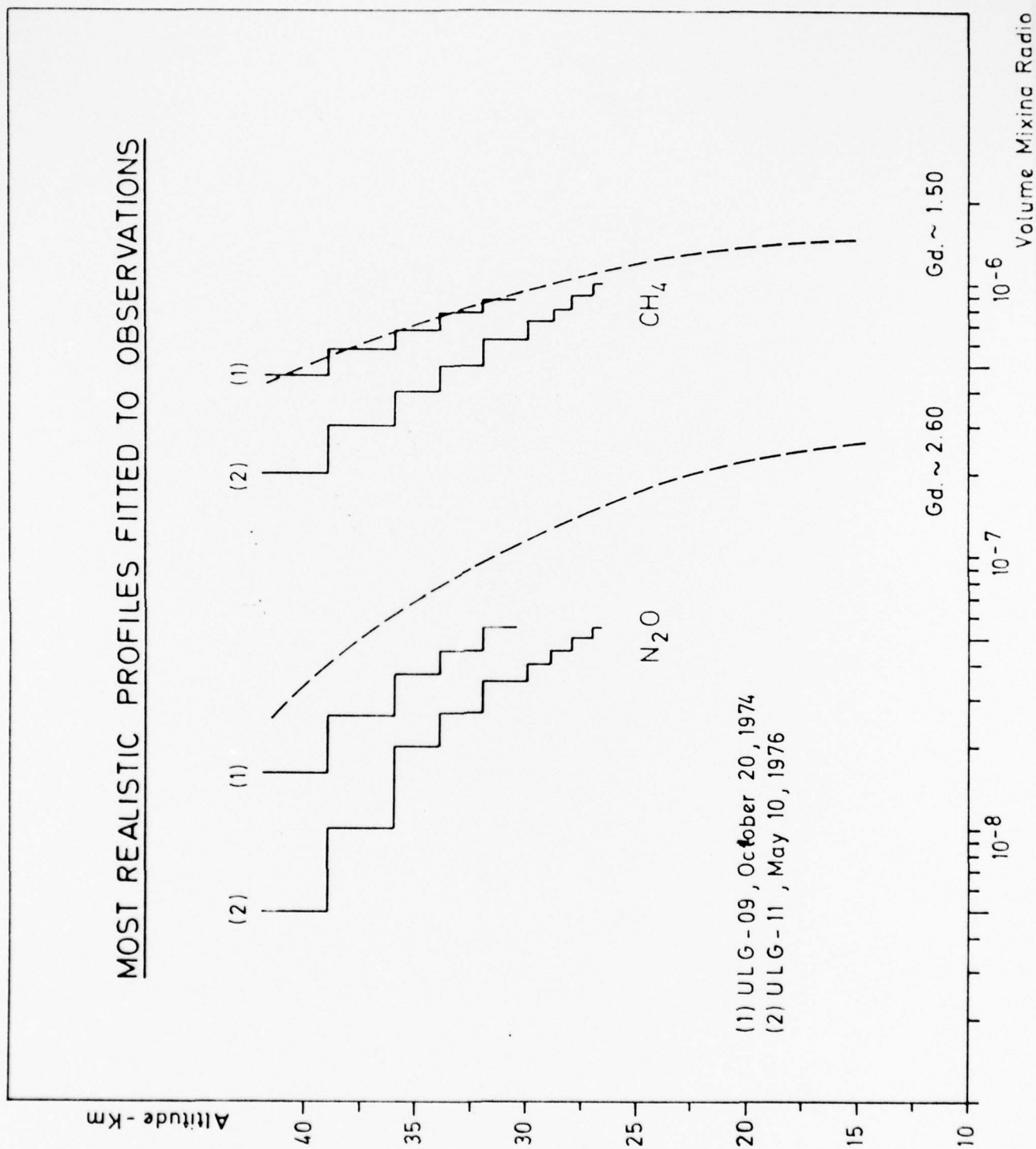
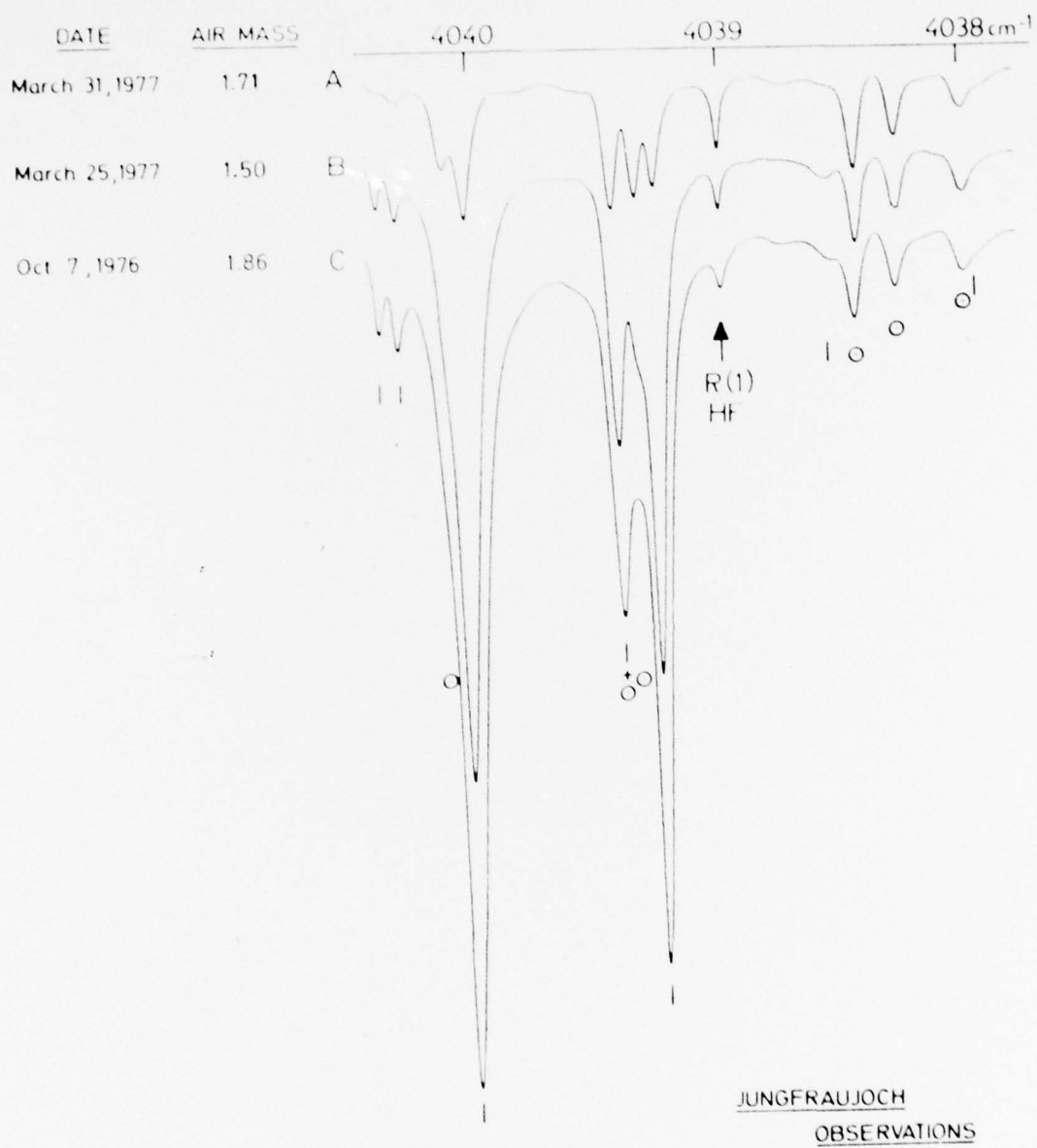


FIGURE 2



Zero A

FIGURE 3

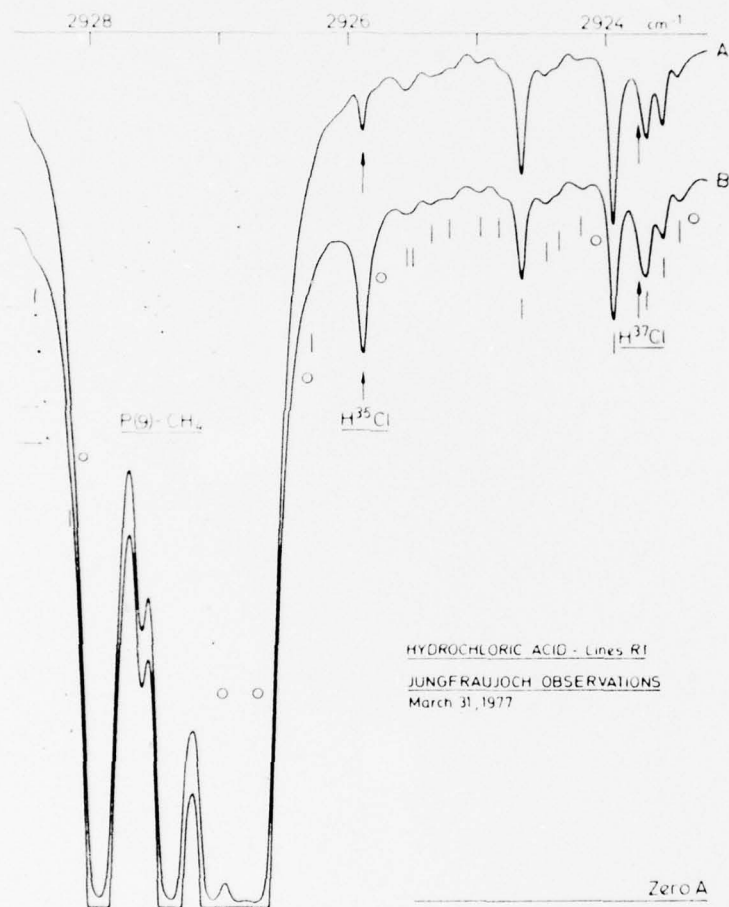


FIGURE 4

Important variations have been noticed among the March observations ; these seem to be due to transport of air from the polar stratosphere, towards lower latitudes and lower altitudes (this is confirmed by similar variations observed on the ozone profiles recorded by the Swiss Meteorological Office).

For HCl :

| | | | | |
|------------------------|--|---|---|-------------------|
| Jungfrauoch (April 51) | average above station : $(1.4 \pm .4) \times 10^{-10}$ ppv | | | |
| (March 77) | " | " | " | $(2.3 \pm .3)$ |
| (July 77) | " | " | " | $(2.2 \pm .3)$ |
| (October 77) | " | " | " | $(2.3 \pm .3)$ |
| (March 78) | " | " | " | $(2.6 \pm .4)$ |
| Kitt Peak (Winter 69) | " | " | " | $(1.5 \pm .4)$ |
| (May 77) | " | " | " | $(1.1 \pm .3) ??$ |

The Kitt Peak 1977 result is surprisingly low and deserves further investigations and measurements ; this is already planed for 1978.

Figures 5 and 6 represent, graphically, the results of the previous tables.

Figure 7 shows a typical spectrum, between 3059 and 3064 cm^{-1} , taken at the Jungfrauoch. The arrows indicate the positions of absorption lines belonging to the ν_4 band of methyl chloride, as calculated by J. S. Margolis and R. A. Toth from the J. P. L., Pasadena, Calif. ; many of these coincide with weak absorption features present on the observed spectra ; till now, the best value deduced for the column density of CH_3Cl above 3580 m altitude is $(1.08 \pm .4) \times 10^{-9}$ ppv.

HF - GROUND MEASUREMENTS

Reduced to Avg. M. Ratio above 20km.

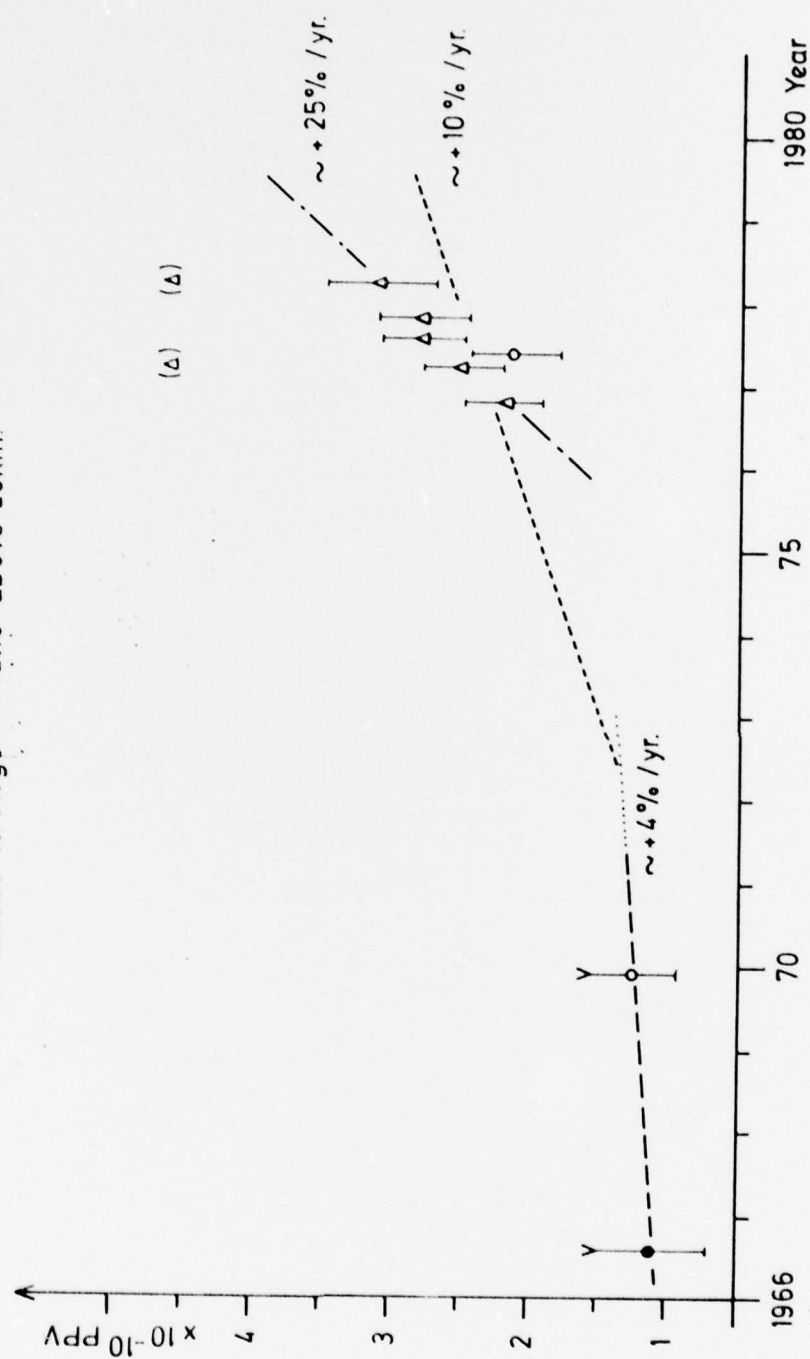


FIGURE 5

HCl - GROUND MEASUREMENTS

Average Mixing Ratio above Station

Δ Jungfraujoeh

o Kitt Peak

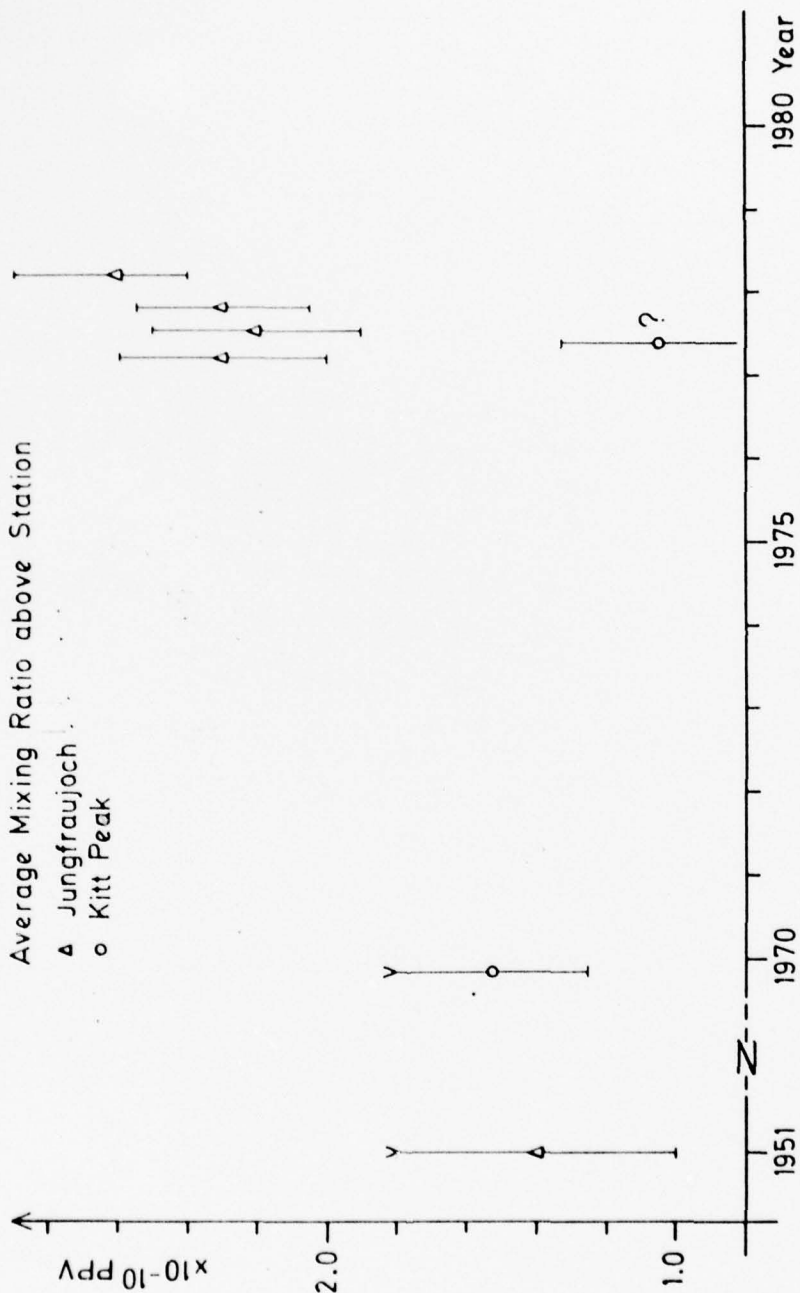
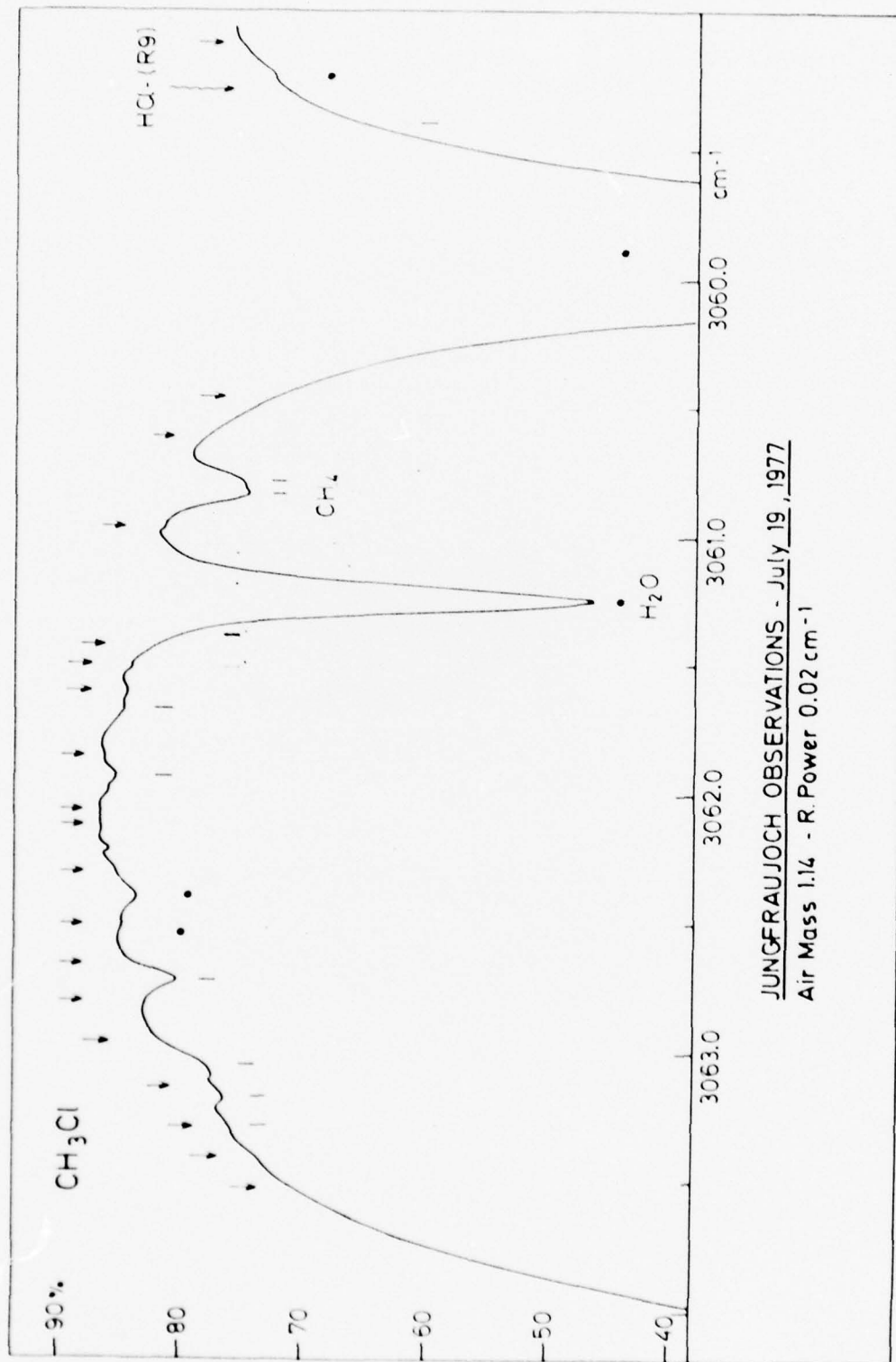


FIGURE 6



JUNGFRAUJOCH OBSERVATIONS - July 19, 1977
 Air Mass 1.14 - R. Power 0.02 cm⁻¹

FIGURE 7

- b. - The observations made during 1977, in the λ 3600- λ 4006 Å and λ 7489- λ 8000 Å are intended to the "Photometric Atlas of the Solar Spectrum, λ 3000 to λ 10,000", published by L. Delbouille, C. Roland and L. Neven ; they will be distributed to the subscribers to that Atlas in a near future.
- c. - The Fourier transform interferometer has been used in 1977, for obtaining solar spectra between 1.1 and 2.5 microns, at the center ($\cos \Theta = 1$) and at the limb ($\cos \Theta = 0.3$) of the solar disk. In particular, the first overtone bands of CO have been observed and they have allowed to deduce indications on the velocity field in the upper layers of the solar photosphere. Reference wavelengths for the 1 to 2.5 microns region have also been deduced with great accuracy (0.001 cm^{-1}) as well as spectroscopic parameters for absorption lines belonging to the $2\nu_3$ band of CH_4 and the following bands of CO_2 : $5\nu_2 + \nu_3 - \nu_2$; $\nu_1 + 4\nu_2 + \nu_3$ and $2\nu_1 + 2\nu_2 + \nu_3$. Most of the spectra obtained with this interferometer have been worked out by Dr. R. Malbrouck and presented in his Ph. D. -thesis, entitled : "Spectroscopie à très haute résolution par transformation de Fourier - Application à l'étude du spectre solaire" [3] .
- Results on telluric CH_4 are summarized in his article reproduced in Appendix B [4] .

In Appendix A, we reproduce a preprint of an article related to another aim of our infrared observational program, i. e. , the observation and identification of weak photospheric lines in the near and middle infrared [5] ; more specific investigations related to that aim have been undertaken by Dr. E. Biémont and presented in his Ph. D. -thesis [6] .

Remark : The 1977 Kitt Peak observations mentioned here have been made by Prof. L. Delbouille and Dr. C. Roland in collaboration with Dr. J. Brault and Mr. L. Testerman from the Kitt Peak National Observatory, Tucson, Arizona, U.S.A.

III. - ACKNOWLEDGMENTS

The efforts reported here have been possible, thanks to financial supports from :

- the belgian government, through the University of Liège ;
 - the Fonds National de la Recherche Scientifique ;
 - the Manufacturing Chemists Association.
-

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Paper Ib 5 - 21st Liège International Astrophysical Symposium
June 21-23, 1977.

APPENDIX A

A SEARCH FOR FAINT MOLECULAR LINES IN THE SOLAR PHOTOSPHERIC
SPECTRUM

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ABSTRACT

On the basis of theoretical predictions concerning the detectability of C_2 , OH and CO bands in the photospheric spectrum, high resolution infrared solar spectra, obtained by balloon, have been investigated. The experimental results deduced for C_2 , in the Ballik-Ramsay band, and for $C^{12}O^{16}$ and $C^{13}O^{16}$, in the fundamental vibration-rotation bands, tend to confirm the calculations.

Results will also be reported from a search for bands of $HS(A^2\Sigma^+-X^2\Pi_i)$, $NH(c^1\Pi-a^1\Delta)$, $CH^+(b^3\Sigma-a^3\Pi)$, $MgH^+(A^1\Sigma-X^1\Sigma)$, $OH^+(A^3\Pi_i-X^3\Sigma^-)$, $HS^+(A^3\Pi_i-X^3\Sigma^-)$, $AlH^+(A^2\Pi-X^2\Sigma)$ and $PH(A^3\Pi_i-X^3\Sigma)$ on near ultraviolet solar spectra recorded at the International Scientific Station, Jungfrauoch (Switzerland).

1. THEORETICAL PREDICTIONS

Calculations of the intensities of a large number of molecular bands which might be present in the solar photospheric spectrum, above 300 nm, have been undertaken, using recent solar photospheric models. From these calculations, we retained a few new systems (see table 1) to be searched for. Table 1 shows that most of the bands retained fall in the near ultraviolet region, whereas three other systems occur in the infrared. Some of these molecules, such as NH, OH, CN, CO, C₂ and CH⁺, present other observed systems in the photospheric spectrum.

The predicted equivalent widths for the strongest lines of each band range between some tenths and a few mÅ for the ultraviolet transitions. The central depths of the infrared lines range between 1 and 6 %.

Large uncertainties, due to the molecular data, do affect the theoretical intensities of some bands :

- (i) lack of experimental $f_{v,v'}$ -values for the systems of MgH⁺ and PH (in such cases, we adopted the values given in Table 1);
- (ii) large uncertainties for the dissociation energy of AlH⁺ and PH;
- (iii) the excitation energy of the lower electronic a state of CH⁺ is poorly known.

2. SEARCH IN THE SOLAR PHOTOSPHERIC SPECTRUM

Using all the available laboratory analysis, we searched for the strongest rotational lines of each of the bands reported in Table 1 on recent high resolution solar tracings obtained during a balloon flight by ~~one~~ of us (Zander, 1976) in the infrared region and by L. Delbouille, L. Neven and G. Roland at the International Scientific Station (Jungraujoch, Switzerland) in the near ultra-violet spectrum.

Our final results are as follows.

C¹³O¹⁶

We have been able to measure, for the first time, lines of the fundamental bands of C¹³O¹⁶ in the solar photospheric spectrum around 5.45 μm (1820-1840 cm^{-1}). Line positions have been computed with the molecular constants of Mantz et al. (1975) for C¹²O¹⁶ and of Todd et al. (1976) for C¹³O¹⁶. From a comparison between weak lines of C¹²O¹⁶ and C¹³O¹⁶, the ratio C¹²/C¹³ has been derived. Figure 1 shows the representative points of the selected lines of the two isotopic species in a curve of growth diagram (linear part). This analysis gives C¹²/C¹³ = 140 \pm 40. While this photospheric result is undoubtedly larger than the terrestrial ratio, C¹²/C¹³ = 90, it agrees however very well with values obtained from recent high resolution studies of the CH, A-X band, around 4300 \AA (Lambert and Mallia, 1968 : 150 \pm 30; Arnould, 1968 : 140 \pm 30; Iversen, 1976 : >90). We are thus in disagreement with results derived from studies of infrared sunspot spectra by Hall et al. (1972) and Hall (1973) who found values in excellent agreement with the terrestrial ratio.

One remark has to be made. The fundamental vibronic bands of CO around $5.4\mu\text{m}$ are certainly the bands to be used for determining the ratio $\text{C}^{13}/\text{C}^{12}$ in the photosphere : the isotopic lines are well separated and they fall in a relatively clear and unperturbed region. On very high resolution spectra with optimum signal to noise ratio (soon available), the $\text{C}^{12}/\text{C}^{13}$ ratio should be derived with a very high accuracy.

C₂ - Ballik-Ramsay

For the first time, lines of the Ballik-Ramsay system of C₂ (0-0 band, around $1.8\mu\text{m}$) have been found in the solar photospheric spectrum recorded by balloon. The observed intensities are in good agreement with predicted intensities based on Cooper and Nicholls' (1975) f-value.

OH - fundamental vibronic band

Preliminary measurements show that many lines of the R-branch near $2.6\mu\text{m}$ seem to be present, as suggested by Migeotte et al. (1957). Interferences with telluric CO₂ and H₂O lines render their identifications somewhat difficult.

Near ultraviolet systems

In this spectral region, the presence of a large number of strong atomic and molecular lines renders any detection of weak line very difficult. Our main purpose was limited to check that none among the lines searched for was really absent from the solar spectrum, owing to the numerous blends. This analysis was made on unpublished solar tracings obtained at the International Scientific Station of the Jungfraujoch (Switzerland) by L. Delbouille, L. Neven and G. Roland.

For each band, only a very small number of lines were rather clean and measurable. Nevertheless, from these few lines, we cannot conclude about the presence of the hereabove mentioned

transitions. All we can say is that, in most of the cases, the smallest observed intensities are of the same order than the predicted ones.

We must also point out that any solar identification strongly depends upon the accuracy of the laboratory line positions. In a first search for OH^+ , made a few years ago by one of us (Sauval, 1972), we concluded to the absence of this molecule from the solar disk spectrum owing to the absence of some of the strongest lines. The present study is simply based on a recent and accurate laboratory analysis by Merer et al. (1975) which disagrees with the older one by Loomis and Brandt (1936).

To conclude with this near ultraviolet search, we shall say that all the bands analyzed are surely not absent and are very probably present in the solar disk spectrum. Laboratory works (line positions, f -values and dissociation energies) are needed to check these identifications.

ACKNOWLEDGMENTS

We thank Drs. L. Delbouille, L. Neven and G. Roland for putting their unpublished solar tracings at our disposal. The help of L. Delbouille and G. Roland during the balloon campaign is also acknowledged. One of us (E.B.) thanks the Fonds de la Recherche Fondamentale Collective for supporting him during this study.

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TABLE 1

Molecular transitions retained for this search

| Molecules | Transition | ΔD_0 (eV) (1) | (v, v'') | $\lambda_{v', v''}$ (Å) | $f_{v', v''}$ (2) | N_{\max} (3) | W_{λ}^{\max} (mÅ) (4) |
|------------------|---------------------------|--------------------------|------------|-------------------------|------------------------|-------------------|----------------------------------|
| N_2H^+ | A $1\Sigma-X$ 1Σ | 0.2 | (0, 2) | 3083 | $(5 \cdot 10^{-3})$ | 15 | 2 |
| NH | c $1\Pi-a$ $1\Delta(x)$ | 0.2 | (0, 0) | 3253 | $3 \cdot 10^{-3}(E)$ | 9 | 0.7 |
| HS | A $2\Sigma^+-X$ $2\Pi_i$ | 0.1 | (0, 0) | 3237 | $1.6 \cdot 10^{-3}(T)$ | 13 | 2 |
| HS ⁺ | A $3\Pi_i-X$ $3\Sigma^-$ | 0.2 | (0, 0) | 3310 | $3.1 \cdot 10^{-3}(E)$ | 13 | 0.3 |
| PH | A $3\Pi_i-X$ 3Σ | 0.4 | (0, 0) | 3420 | (10^{-2}) | 13 | 0.3 |
| CH ⁺ | b $3\Sigma-a$ $3\Pi(x)$ | <0.1 | (0, 0) | 3490 | $2 \cdot 10^{-3}(E)$ | 11 | 0.2 |
| OH ⁺ | A $3\Pi_i-X$ $3\Sigma^-$ | 0.2 | (0, 0) | 3565 | $4.6 \cdot 10^{-3}(E)$ | 9 | 0.6 |
| AlH ⁺ | A $2\Pi-X$ 2Σ | 0.5 | (0, 0) | 3612 | $10^{-2}(T)$ | 16 | 1 |
| C_2 | b $3\Sigma_g^-a$ $3\Pi_u$ | <0.1 | (0, 0) | 17675 | $10^{-3}(E)$ | 30 | 2 |
| OH | X $2\Pi, \Delta v=1$ | <0.1 | - | 26000 | $5 \cdot 10^{-6}(E)$ | 9 | 40 |
| $C^{13}O^{16}$ | X $1\Sigma, \Delta v=1$ | <0.1 | - | 54500 | $10^{-5}(E)$ | 28 | 60 |

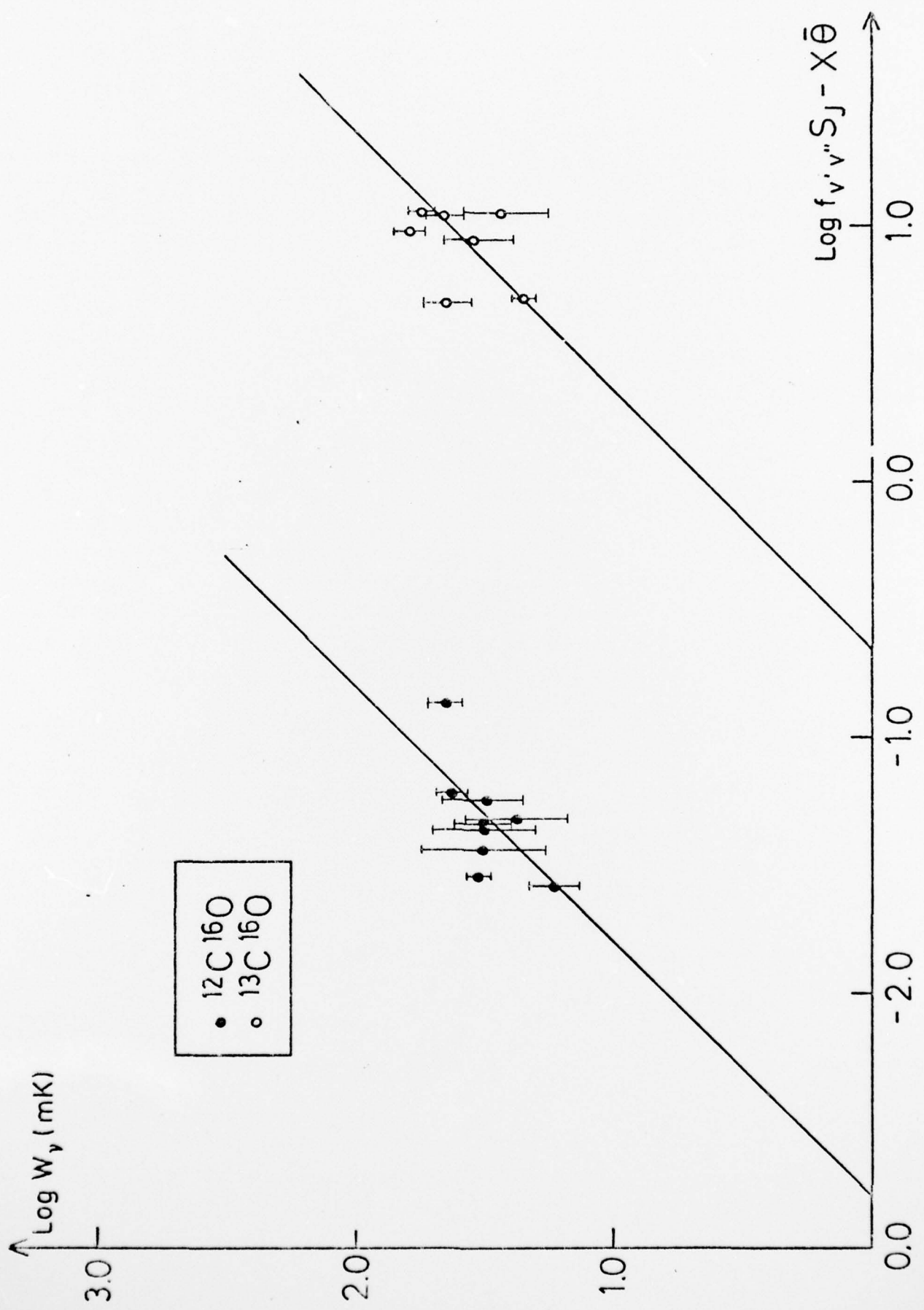
TABLE 1. (continued)

- (1) uncertainty of the dissociation energy, D_o , in eV;
- (2) E : experimental result; T : theoretical result;
- (3) N_{\max} : rotational quantum number for which maximum intensity is reached in the solar spectrum;
- (4) W_{λ}^{\max} : solar equivalent width corresponding to N_{\max} , in mÅ;
- (x) For NH, we adopted $T_e = 1.56$ eV; for CH^+ , $T_e = 1.0$ eV.

FIGURE CAPTION

Figure 1.— Linear parts of the curves of growth for $C^{12}O^{16}$ (filled circles) and $C^{13}O^{16}$ (open circles). The ordinate give the logarithm of the equivalent width, in mK; the abscissa gives the quantity $\log f_{v,v''} S_J - \chi \bar{\theta}$, where S_J is the ~~oscillator~~ ^{Hönl}-London factor, χ , the excitation energy and $\bar{\theta} = \frac{5040}{T}$ is a reciprocal mean temperature of the layer where the lines are formed (we took $\bar{\theta} = 1.05$).

Fig. 1. Unravelling of the ...



Extrait du *Bulletin de l'Académie royale de Belgique* (Classe des Sciences)
Séance du samedi 5 novembre 1977.

**Abondance du méthane tellurique déterminée
à partir de l'analyse de la bande $2\nu_3$**

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Résumé. — Deux spectres solaires à très haute résolution ont été enregistrés par la méthode de spectroscopie par transformation de Fourier à la Station Scientifique Internationale du Jungfraujoch, dans la région de la bande $2\nu_3$ du méthane. Les positions précises des raies de cette bande ainsi que leurs demi-largeurs ont été déterminées. Le taux de mélange moyen du méthane dans l'atmosphère terrestre, déduit de l'analyse de nombreux multiplets de la bande $2\nu_3$, est de $1.48 \cdot 10^{-6}$.

Abstract. — Two very high resolution solar spectra have been recorded by Fourier spectroscopy at the International Scientific Station of the Jungfraujoch, in the region of the $2\nu_3$ band of methane. Accurate positions of the lines of these band and their half-width have been determined. The mean mixing ratio of methane in the earth atmosphere deduced from the analysis of numerous multiplets of the $2\nu_3$ band is $1.48 \cdot 10^{-6}$.

1. INTRODUCTION

Le méthane a été considéré comme un constituant mineur de l'atmosphère terrestre depuis que la bande fondamentale ν_3 à 3,3 microns a été identifiée pour la première fois par Migeotte dans le spectre solaire [1-3]. Depuis lors, un certain nombre de déterminations de l'abondance du méthane tellurique ont été effectuées sur la base d'observations spectroscopiques de la bande $2\nu_3$ située à 1,6 microns.

En 1949, McMath *et al.* [4] ont déduit, à partir d'un spectre solaire enregistré dans cette région, une température rotationnelle de 236°K et

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un taux de mélange de l'ordre de $1,2 \cdot 10^{-6}$. Deux ans plus tard, Gold-b[1] détermine des températures rotationnelles de 229°K et 227°K et des abondances de CH_4 de 1,19 et 0,967 cm-atm, à partir d'observations de la bande $2\nu_3$ effectuées à Lake Angelus (296 m) et au Mont Wilson (1742 m).

En 1962, Nielsen et Migeotte [6] effectuèrent des mesures similaires à la Station Scientifique Internationale du Jungfraujoch (Suisse, Alt: 3850 m) et obtinrent $T_{\text{rot}} = 206^\circ\text{K}$ et une abondance de 0,777 cm-atm.

Plus récemment, Margolis *et al.* [7] ont redéterminé la température rotationnelle pour la bande $2\nu_3$ ainsi que l'abondance du méthane tellurique à partir d'un spectre solaire enregistré à l'aide d'un interféromètre de type Connes (limite de résolution $0,10 \text{ cm}^{-1}$) au McDonald Observatory et ont obtenu $T_{\text{rot}} = (243,4 \pm 5,3)^\circ\text{K}$ et une abondance de $1,137 \pm 0,117 \text{ cm-atm}$ (STP) ce qui correspond à un taux de mélange de $1,55 \cdot 10^{-6}$.

Depuis l'installation, en juin 1974, d'un interféromètre de type Connes à la Station Scientifique du Jungfraujoch, des spectres solaires à très haute résolution ont été enregistrés dans les fenêtres atmosphériques situées dans le domaine de sensibilité des cellules au sulfure de plomb (1 à 3,5 microns). Deux de ces enregistrements effectués respectivement le 27 février 1975 et le 21 mai 1975 dans la région de la bande $2\nu_3$ du CH_4 , aux environs de 1,6 microns, ont permis de remesurer l'abondance intégrée du méthane tellurique. Dans la réduction des observations, les données de laboratoire les plus récentes concernant les intensités standards et le coefficient de température ont été utilisées.

D'autre part, la très haute résolution des enregistrements a permis non seulement de préciser la position des composantes non perturbées des multiplets de la bande $2\nu_3$ mais également de mesurer les demi-largeurs des raies qui sont des paramètres fondamentaux dans les problèmes d'absorption atmosphérique.

2. OBSERVATIONS

Les spectres solaires ont été obtenus à l'aide d'un interféromètre de type Connes de 1 mètre de différence de marche, ce qui permet d'atteindre une limite de résolution de $0,01 \text{ cm}^{-1}$ (avec l'apodisation la plus forte). Cet instrument est installé au foyer coudé du télescope de 76 cm de la Station Scientifique Internationale du Jungfraujoch. Il

fonctionne actuellement suivant la méthode du pas à pas avec intégration des échantillons de l'interférogramme pendant les périodes de stabilisation de la différence de marche aux positions d'échantillonnage, avec une précision de quelques Angströms.

Les domaines spectraux sont isolés à l'aide d'un monochromateur à dispersion soustractive placé entre le télescope et l'interféromètre [8].

Les caractéristiques essentielles des enregistrements solaires qui ont servi de base à cette étude sont reprises dans le tableau I.

TABLEAU I. — Caractéristiques des enregistrements solaires.

| | | |
|---------------------------|----------------------------|----------------------------|
| Date | 27-02-1975 | 21-05-1975 |
| Limite de résolution | 0,017 cm^{-1} | 0,015 cm^{-1} |
| Durée des enregistrements | 2 h 30 | 2 h 15 |
| Massé d'air | 1,76 à 1,98 | 1,13 à 1,16 |
| Domaine spectral analysé | 5820-6150 cm^{-1} | 5900-6400 cm^{-1} |

Le rapport signal/bruit des spectres est supérieur à 100:1. Ces spectres solaires couvrent toute la région de la bande $2\nu_3$ du méthane depuis P(10) jusqu'à R(9).

La figure 1 reproduit le tracé photométrique correspondant à l'enregistrement effectué le 21-05-1975, dans la région de la branche Q de la bande $2\nu_3$. Les divers multiplets de cette branche sont indiqués sur la figure, les autres raies sont d'origine solaire et sont caractérisées par le symbole \odot .

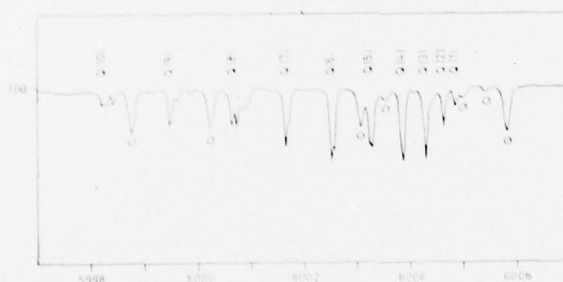


FIG. 1. — Le spectre solaire dans la région de la branche Q de la bande $2\nu_3$ du CH_4 atmosphérique.

3. RÉDUCTION DES DONNÉES

En supposant un profil de Lorentz, la transmission Tr dans le voisinage d'une raie d'absorption isolée est donnée par l'expression:

$$Tr(\sigma) = \exp(-k(\sigma)l) \quad (1)$$

où l est la longueur du trajet d'absorption et σ le nombre d'ondes. Le coefficient d'absorption $k(\sigma)$ est donné par la relation:

$$k(\sigma) = \frac{S}{H(\sigma - \sigma_0)^2 + \gamma_L^2} \quad (2)$$

dans laquelle S est l'intensité de la raie, γ_L sa demi-largeur et σ_0 le nombre d'ondes central de la raie.

La demi-largeur γ_L est fonction de la pression P et de la température T et se met sous la forme:

$$\gamma_L = \gamma_0 \frac{P}{P_0} \left(\frac{T_0}{T} \right)^n \quad (3)$$

où n représente le coefficient de température.

Les grandeurs indicées par un zéro représentent des valeurs déterminées dans des conditions standard de température et de pression.

Le coefficient de température adopté dans ce travail est celui qui a été publié par Darnton et Margolis [9] et qui a été déterminé en laboratoire, à partir des raies $P(2)$, $R(0)$, $R(1)$ et $R(2)$ de la bande $2\nu_3$.

L'intensité S est définie par la relation:

$$S = \int k(\sigma) d\sigma \quad (4)$$

Elle est seulement fonction de la température et a pour expression:

$$S(T) = S_0(T/T_0)^{1/2} \exp(-BJ(J+1)(hc/k)((T_0 - T)/T_0T)) \quad (5)$$

où B est la constante rotationnelle ($5,24 \text{ cm}^{-1}$) et J le nombre quantique de rotation.

La largeur équivalente d'une raie est définie par la relation:

$$W = \int |1 - (I/I_0)_\sigma| d\sigma \quad (6)$$

L'intensité d'une raie, sa largeur équivalente et sa demi-largeur vérifient l'équation de Ladenburg-Reiche:

$$W = 2H\gamma_L F(x) \quad (7)$$

où $F(x) = xe^{-x} |J_0(ix) - iJ_1(ix)|$ dans laquelle $J_0(ix)$ et $J_1(ix)$ représentent des fonctions de Bessel d'ordre 0 et 1. La variable x est liée à la quantité d'absorbant l et s'écrit:

$$x = Sl/2H\gamma_L. \quad (8)$$

4. ANALYSE DES SPECTRES

A. POSITION DES RAIES

Un des grands avantages de la méthode de spectroscopie par transformation de Fourier, par rapport aux méthodes classiques de spectroscopie, est de fournir des mesures très précises des nombres d'ondes des raies spectrales. Théoriquement, la précision que l'on peut atteindre sur la valeur absolue du nombre d'ondes d'une raie présente dans un spectre obtenu par spectroscopie de Fourier n'est limitée que par celle que l'on a sur la valeur absolue de la longueur d'onde de la raie de référence qui sert au contrôle de la différence de marche.

Les positions des raies non perturbées de la bande $2\nu_3$ présentes dans les spectres solaires ont été déterminées à l'aide d'un programme spécial de détection automatique de raies, mis à notre disposition par J. Chauville et J.P. Maillard [9] de l'Observatoire de Meudon.

Le tableau II reproduit les nombres d'ondes, dans le vide, des diverses composantes des multiplets des branches P, Q et R de la bande $2\nu_3$. Les raies sont identifiées selon les notations utilisées par Bobin [10]. La notation * indique que les composantes correspondantes ne sont pas résolues sur les spectres solaires.

B. DEMI-LARGEUR DES RAIES

La demi-largeur des raies est un paramètre fondamental dans les problèmes d'absorption atmosphérique.

En ce qui concerne la bande $2\nu_3$ de CH_4 , Darnton et Margolis [11] ont étudié l'influence de la température sur la demi-largeur des raies.

TABLEAU II. — Nombre d'ondes des raies de la bande $2\nu_3$ de CH_4 .

| Branche P | | Branche Q | | Branche R | |
|-------------|------------------------|-------------|------------------------|-------------|------------------------|
| Raie | σ (cm $^{-1}$) | Raie | σ (cm $^{-1}$) | Raie | σ (cm $^{-1}$) |
| P1 F $_1$ | 5994.147 | Q1 F $_1$ | 6004.866 | R0 A $_1$ | 6015.667 |
| P2 F $_2$ | 5983.192 | Q2 F $_2$ | 6004.651 | R1 F $_1$ | 6026.233 |
| P3 A $_2$ * | 5972.121 | Q3 A $_2$ | 6004.309 | R2 F $_2$ | 6036.660 |
| F $_1$ * | | Q4 F $_1$ | 6003.884 | R3 A $_2$ | 6046.959 |
| P4 F $_2$ * | 5960.881 | Q5 F $_1$ | 6003.262 | R4 F $_2$ * | 6057.094 |
| F $_1$ * | | F $_2$ | 6003.305 | F $_1$ * | |
| P5 F $_1$ | 5949.614 | Q6 A $_1$ | 6002.532 | R5 F $_1$ | 6067.151 |
| F $_2$ * | 5949.544 | A $_2$ | 6002.589 | F $_2$ | 6067.097 |
| F $_1$ * | | Q7 F $_1$ | 6001.674 | R6 A $_1$ | 6077.049 |
| P6 A $_1$ * | 5938.181 | F $_2$ | 6001.737 | A $_2$ | 6076.950 |
| F $_2$ * | | Q8 F $_2$ * | 6000.659 | R7 F $_1$ | 6086.789 |
| A $_2$ | 5938.092 | F $_1$ * | | A $_2$ | 6086.755 |
| P7 F $_1$ | 5926.649 | Q9 A $_2$ * | 5999.498 | F $_2$ | 6086.633 |
| F $_2$ * | 5926.590 | A $_1$ * | | R8 F $_2$ | 6096.505 |
| A $_2$ * | | F $_2$ | 5999.631 | F $_1$ * | 6096.381 |
| F $_2$ | 5926.479 | | | F $_2$ * | |
| P8 F $_2$ | 5915.051 | Q10 F $_1$ | 5998.233 | F $_1$ | 6096.176 |
| F $_2$ | 5914.918 | A $_1$ * | 5998.318 | R9 F $_2$ | 6106.251 |
| F $_1$ | 5914.765 | F $_1$ * | | | |
| P9 A $_2$ | 5903.362 | A $_2$ | 5998.442 | F $_1$ | 6106.050 |
| F $_1$ * | 5903.308 | | | F $_2$ | 6105.628 |
| A $_1$ * | | | | | |
| F $_1$ | 5903.172 | | | | |
| F $_2$ * | 5902.959 | | | | |

Malheureusement, ces mesures ne concernent que 4 raies de la bande: les raies P(2), R(0), R(1) et R(2).

La très haute résolution de nos enregistrements solaires a permis de mesurer la demi-largeur de toutes les raies non perturbées de la bande $2\nu_3$, dans nos conditions d'observation (pression au niveau du sol à l'altitude du Jungfraujoch de 640 mb).

La figure 2 reproduit les demi-largeurs (en mK) mesurées sur les spectres et corrigées du profil instrumental, en fonction du nombre quantique rotationnel J . Les mesures effectuées sur la branche R sont représentées par des points, celles obtenues sur la branche Q, par des cercles et celles correspondant à la branche P, par des triangles.

Cette figure fait apparaître une dépendance des demi-largeurs des raies en fonction de J . Il faut cependant signaler que des effets de saturation peuvent intervenir pour les raies les plus intenses des branches de la bande et plus particulièrement au voisinage de $J = 5$. La

valeur de γ correspondant à la raie R(0) est de 0.030 cm^{-1} ; elle est en bon accord avec celle que l'on déterminerait en prenant la valeur publiée par Darnton et Margolis [11] et en l'extrapolant pour nos conditions d'observation.

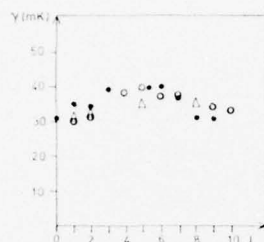


FIG. 2. — Demi-largeurs des raies de la bande $2\nu_3$ du méthane atmosphérique déterminées à partir des spectres solaires.

5. ABONDANCE DU CH_4 TELLURIQUE AU-DESSUS DU JUNGFRAUJOCH

L'intensité S des raies a été calculée à l'aide de la relation [5] à partir des intensités standard S_0 publiées par Margolis [12] en se basant sur les fréquences spectrales données par Bobin [10]. La température rotationnelle adoptée est celle qui a été déterminée par Margolis *et al.* [7].

Les largeurs équivalentes de 15 multiplets appartenant essentiellement aux branches R et Q ont été mesurées par planimétrie et corrigées de la contribution due aux ailes.

La quantité de CH_4 tellurique, au-dessus de l'altitude du Jungfraujoch déterminée à l'aide du formalisme décrit au paragraphe 3 est donnée dans le tableau III pour les deux observations.

TABLEAU III. Quantité de CH_4 tellurique au-dessus de l'altitude du Jungfraujoch (cm-atm)_{STP}.

| Date | 27-02-1975 | 21-05-1975 |
|----------|-------------------|-------------------|
| Quantité | 0.781 ± 0.079 | 0.762 ± 0.026 |

Si l'on tient compte de la hauteur d'échelle de l'atmosphère à l'altitude du Jungfraujoch, le taux de mélange moyen de CH_4 correspondant est de $1.48 \cdot 10^{-6}$, légèrement inférieur à celui obtenu par Margolis *et al.* [7].

CONCLUSIONS

L'abondance du méthane tellurique a été déterminée à partir de deux spectres solaires à très haute résolution obtenus par spectroscopie de Fourier à la Station Scientifique Internationale du Jungfraujoch (3880 m).

De nombreux multiplets de la bande $2\nu_3$ ont été analysés, ce qui a permis de fournir des positions précises des raies et de déterminer, grâce à la très haute résolution des enregistrements, leur demi-largeur.

Le choix de la bande $2\nu_3$ est particulièrement bien indiqué pour des dosages atmosphériques étant donné que la majorité des raies de cette bande se situent sur la partie linéaire de la courbe de croissance et que l'on possède des données de laboratoire précises concernant les intensités standard et le coefficient de température.

Les deux spectres solaires analysés ne montrent pas de variation importante de la quantité de CH_4 tellurique. Notons, en particulier, que la moyenne de nos mesures est en excellent accord avec celle obtenue en 1962 par Nielsen et Migeotte. Nous nous proposons de répéter régulièrement ces mesures spectroscopiques de l'abondance du CH_4 en enregistrant de façon régulière le spectre solaire dans la région de la bande $2\nu_3$, depuis le Jungfraujoch, afin de détecter d'éventuelles variations de la concentration du méthane dans l'atmosphère terrestre.

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